

#### Phase space transport. II

#### **Shane Ross**

Control and Dynamical Systems, Caltech http://www.cds.caltech.edu/~shane/

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# Motivation

Apply transport calculations to asteroid pairs to calculate, e.g., capture & escape rates.

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Dactyl in orbit about Ida, discovered in 1994 during the Galileo mission.

# Motivation

- □ Movie of simple model of Ida & Dactyl
- □ Small hard sphere around rotating elliptical asteroid
- Nonmerging collisions modeled as bounces (work with E. Kanso)

# In This Talk...

#### Part I. Restricted F2BP phase space

- □ Dependence on energy
- □ Tube + lobe dynamics
- Modeling collisions
- Part II. Transport using set oriented methods
  - □ Transport problem described
  - Two computational techniques
    - (a) Invariant manifolds
    - (b) Almost-invariant sets
  - Extensions and future work

### Part I

#### **Restricted Full Two Body Problem**

- $\Box$  Consider two masses:  $m_1$  (sphere) &  $m_2$  (ellipse)
- $\begin{array}{ccc} \frac{m_1}{m_2} \rightarrow 0 & m_1 & m_2 \\ \text{Particle around asteroid} & & & & \\ \text{restricted F2BP (RF2BP)} & & & & & & \\ \end{array}$
- Restricted (as in restricted 3-body problem) simple case exhibits the basic capture, ejection, collision dynamics (see Koon, Marsden, Ross, Lo, Scheeres [2003])
- □ Can include bouncing, sticking, ... (see Kanso [2003])

## **Restricted F2BP**

Point mass P moving in the x-y plane under the gravitational field of a uniformly rotating elliptical body m, without affecting its uniform rotation.



The rotating (x-y) and inertial (X-Y) frames.

### **Restricted F2BP**

Equations of motion relative to a rotating Cartesian coordinate frame and appropriately normalized:

$$\ddot{x} - 2\dot{y} = -\frac{\partial U}{\partial x}$$
 and  $\ddot{y} + 2\dot{x} = -\frac{\partial U}{\partial y}$ ,

where

$$U(x,y) = -\frac{1}{r} - \frac{1}{2}r^2 - \frac{3C_{22}\left(x^2 - y^2\right)}{r^5},$$

and

$$r = \sqrt{x^2 + y^2}.$$

**Energy integral (Jacobi integral):**  $E = \frac{1}{2} \left( \dot{x}^2 + \dot{y}^2 \right) + U(x, y).$ 

### **Restricted F2BP**

 $\Box$  Gravity field coefficient  $C_{22}$ , the ellipticity,

$$C_{22} = \frac{1}{20}(1 - \beta^2),$$

varies between 0 and 0.05.

 $\Box$  e.g., Ida:  $\beta \approx 0.43, C_{22} \approx 0.04$ 



# Phase Space Structure

Energy indicates type of global dynamics.

Is movement between the exterior and asteroid realms possible?



# Phase Space Structure

#### Multi-scale dynamics

Coarse level : tube dynamics between realms
 Fine level : lobe dynamics within realms

# Phase Space Structure

Slices of energy surface: Poincaré sections U<sub>i</sub>
 Tube dynamics: evolution between U<sub>i</sub>
 Lobe dynamics: evolution on U<sub>i</sub>



 $\Box$  Suppose  $E < E_S$ ; energy surface  $\mathcal{M}_E$ 

- □ Asteroid and exterior realms not connected
- $\Box$  Poincaré map in exterior realm: area and orientation preserving map on  $M\subset \mathbb{R}^2$ ,

$$f: M \longrightarrow M$$

where

$$M = \mathcal{M}_E \cap \{x = 0, \dot{x} > 0\}$$

with coordinates  $(y,p_y)$ , equiv.,  $(r,p_r)$ , on M

□ Particles are **ejected** if they lie within lobes enclosed by the stable and unstable manifolds of a hyperbolic fixed point at  $(+\infty, 0)$ —lobes of ejection.





Numerical simulation using MANGEN



Curves can be followed to very high accuracy

# **MANGEN** Description

Simulations use MANGEN (Coulliete & Lekien)
 Adaptive conditioning of curves based on curvature.



#### $\Box$ Suppose $E > E_S$

- Exterior and asteroid realms **connected via tubes**
- In exterior realm, some tubes lead to collision (others lead away from collision)

#### $\Box$ Tube + lobe dynamics =

Alternate fates of collision and ejection are intimately intermingled.

#### □ Tubes leading to collision with asteroid



Position space projection

 ${\rm Motion} \, \, {\rm on} \, \, M$ 

#### Tubes leading to collision with asteroid plus tubes coming from collision, e.g., liberated particles



#### Escape and re-capture.



# **Collision Modeling**

If bouncing is modeled, dynamics is more complicated.
 Upon bouncing, particle moves to new energy surface
 Work in progress with E. Kanso



## Part II

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## Part II

#### Transport using set oriented methods

- $\Box$  Describe transport of phase points on a k-dimensional manifold M
- $\hfill M$  could be, e.g., the ocean surface, an energy shell, or a Poincaré surface-of-section
- $\Box$  We look first at k = 2 for autonomous systems
- Paper: Dellnitz, Junge, Koon, Lekien, Lo, Marsden, Padberg, Preis, R., Thiere [2003]

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 $\Box$  Initially,  $R_i$  is uniformly covered with species  $S_i$ . i.e., Species type indicates where a point was initially.

 $\Box$  Describe the distribution of species  $S_i$  throughout the regions  $R_i$  at any future iterate n > 0.

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# **Transport Quantities**

#### **Quantities of interest:**

 $T_{i,j}(n) \equiv$  the total amount of species  $S_i$  contained in region  $R_j$  immediately after the *n*-th iterate  $= \mu(f^{-n}(R_j) \cap R_i)$ 

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#### Our goal:

Compute the  $T_{i,j}(n)$  up to some  $n_{\max}$ 

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Compare & combine two computational approaches

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# **Computational Approaches**

Compare & combine two computational approaches

- 1) Invariant manifolds of fixed points, lobe dynamics; co-dimension one objects which bound regions, etc.
  MANGEN: Manifold Generation, Lekien etc
- 2) Set oriented methods, almost-invariant sets;
  direct computation of regions
  GAIO: Global Analysis of Invariant Objects, Dellnitz etc

# Particle in 2-Body Field

Example problem: test particles in gravity field of two masses,  $m_1$  and  $m_2$ , in circular orbit, i.e., the planar, circular restricted three-body problem with  $\frac{m_2}{m_1} \approx 10^{-3}$ .

□ Reduce to 2D map via Poincaré surface-of-section



# Particle in 2-Body Field

• Poincaré map  $f: M \to M$  has regular and irregular components. Large connected irregular component, the "chaotic sea."



To understand the transport of points under the Poincaré map *f*, we consider the invariant manifolds of unstable fixed points

Let  $p_i, i = 1, ..., N_p$ , denote a collection of saddle-type hyperbolic fixed points for f.

Local pieces of unstable and stable manifolds







Unstable and stable manifolds in **red** and **green**, resp.

• Intersection of unstable and stable manifolds define boundaries.



• These boundaries divide phase space into regions,  $R_i, i = 1, \ldots, N_R$ 



#### □ Local transport: across a boundary

consider small sets bounded by stable & unstable mfds



- They map from entirely in one region to another under one iteration of f
  - $L_{1,2}(1)$  and  $L_{2,1}(1)$  are called turnstile lobes



 $\square$  MANGEN: evolution of a lobe of species  $S_1$  into  $R_2$ 

insert S1 into R2 movie

**Global transport** between regions  $(T_{i,j}(n))$  is completely described by the dynamical evolution of lobes.



# Set Oriented Methods

#### **Overview**

Partition phase space into loosely coupled regions

$$R_i, i=1,\ldots,N_R,$$

 $\Box$  Probability is small for a point in a region to leave in a short time under f.

□ These **almost-invariant sets** (AIS's) define macroscopic structures preserved by the dynamics.

 $\Box$  The transport,  $T_{i,j}(n)$ , between almost-invariant sets can then be determined.

1) discretize the phase space into boxes; model boxes as the vertices and transitions between boxes as edges of a directed graph



2) use graph partitioning methods to divide the vertices of the graph into an optimal number of parts such that each part is highly coupled within itself and only loosely coupled with other parts



□ 3) by doing so, we can obtain AIS's and analyze transport between them

![](_page_48_Figure_2.jpeg)

#### **Box Formulation**

Create a fine box partition of the phase space
 \$\mathcal{B} = {B\_1, \ldots B\_q}\$, where \$q\$ could be \$10^7\$+
 Consider a (weighted) \$q\$-by-\$q\$ transition matrix, \$P\$, for our dynamical system, where

$$P_{ij} = \frac{\mu(B_i \cap f^{-1}(B_j))}{\mu(B_i)},$$

the transition probability from  $B_i$  to  $B_j$ 

 $\Box P$  is an approximation of our dynamical system via a finite state Markov chain.

#### **Graph Formulation and Partitioning**

 $\Box P$  has a corresponding graph representation where nodes of the graph correspond to boxes  $B_i$ .

![](_page_50_Figure_3.jpeg)

- □ If  $P_{ij} > 0$ , then there is an edge between nodes *i* and *j* in the graph with weight  $P_{ij}$ .
- Partitioning into AIS's becomes a problem of finding a minimal cut of this graph.

![](_page_51_Figure_3.jpeg)

#### □ **AIS's correspond with key dynamical features** More refined methods like MANGEN can pick up details

![](_page_52_Figure_2.jpeg)

The phase space is divided into several invariant and almost-invariant sets.

Using the box formulation and GAIO, the  $T_{i,j}(n)$  can be computed for large n. Agrees with MANGEN result.

![](_page_53_Figure_2.jpeg)

To speed the computation, box refinements are performed where transport related structures, e.g., lobes, are located.

![](_page_54_Figure_2.jpeg)

□ The merging of **statistical** and **geometric** approaches yields a very powerful tool.

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- Example problem: restricted 3-body problem.
- □ Both find the same regions
  - AIS's, statistical features, are identified with regions, geometric features

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 $\Box$  Same values for  $T_{i,j}(n)$  over common time window.

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  - GAIO for coarse picture, transport calculations MANGEN to refine on regions of interest
  - $\Rightarrow$  important for precision navigation

- $\Box$  AIS & lobe dynamics in 3D+, e.g., astronomy, chemistry
- □ AIS for time dependent systems? e.g., ocean dynamics
- □ Numerical algorithms are crucial:
  - GAIO for coarse picture, transport calculations MANGEN to refine on regions of interest
  - $\Rightarrow$  important for precision navigation
- □ Merge techniques into single package:
  - Box formulation, graph algorithms
  - Co-dimension one objects
  - Adaptive conditioning based on curvature

# **Selected References**

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For papers, movies, etc., visit the websites: http://www.cds.caltech.edu/~shane http://www.nast-group.caltech.edu/

![](_page_66_Picture_0.jpeg)

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